AD/A-001 238

AERODYNAMICS OF IN-TUBE LAUNCH PHASE OF ROCKETS (INTERNAL BALLISTICS ANALYSIS AND MINI-COMPUTER PROGRAM DEVELOPMENT)

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Illinois University

Prepared for:

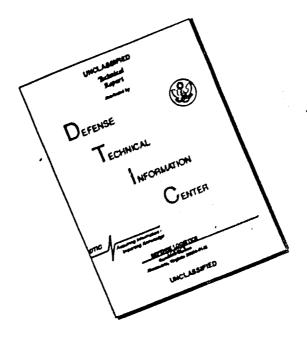
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REPORT DOCUMENTATION PAGE	READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER 2. GOV'T ACCESSION N	O. 3 RECIPIENT'S CATALOG NUMBER
UILU ENG 74-4006	ADIA-001 2.38
4. TITLE (mid Subtitio)	S. TYPE OF REPORT & PERIOD COVERED
AERODYNAMICS OF IN-TUBE LAUNCH PHASE OF ROCKETS	Final, 1 November 1973-30 Jun
(INTERNAL BALLISTICS ANALYSIS AND MINI-COMPUTER DEVELOPMENT)	6. PERFORMING ORG. REPORT NUMBER
7. AUTHOR(e)	8. CONTRACT OR GRANT NUMBER(4)
Helmut H. Korst	DAHC04-74-G-0057
PERFORMING ORGANIZATION NAME AND ADDRESS	10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
Gas Dynamics Laboratory, Department of Mech. & Ind. Eng., University of Ill. at Urbana-Champaig Urbana, IL 61801	Post LRCP-U.S.A.R.O.
1. CONTROLLING OFFICE NAME AND AGORESS	12. REPORT DATE
Department of the Army	
U.S.A.R.O., Box CM, Duke Station	Thirty-Three (33)
Durham NC 27706 14. MONITORING SGENCY NAME & ADDRESS(II dillerent from Controlling Office)	15. SECURITY CLASS. (of this report)
U.S.A. Missile Command	
Redstone Arsenal, AL 35809	
national in Jenui, he Journ	150. DECLASSIFICATION OOWN GRADING SCHEDULE
b. DISTRIBUTION STATEMENT (of this Report)	
Approved for Public Release Distribution Unlimited 2. DISTRIBUTION STATEMENT (of the observed entered in Block 20, if different in	Page 1
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Internal Ballistics Rocket Launch Mini-Computer Programs

20. AUSTRACT (Continue on reverse olds if necessary and identity by block number)

Aerodynamic phenomena affecting the launch phase of in-tube launched rockets are investigated. Of special interest is the motion of the missile relative to the tube before and during exiting as it results from the burning characteristics of the propellant with given grain geometry, structural constraints such as closure failure, and tube wall proximity.

(continued)

20. ABSTRACT (continued)

Theoretical analysis, based on certain simplifying assumptions such as uniform, only pressure dependent grain burning, one-dimensional quasi-steady nozzle flow, and constant flame temperature, results in systems of ordinary non-linear differential equations which have been programmed for mini-computer use.

Quantitative results are obtained which show good agreement with available data and describe the motion of the rocket during launch phase in sufficient detail. Rocket chamber pressure and nozzle thrust are evaluated over the entire duration of powered flight.

An experimental program in support of the continuing effort has been initiated by adding high-pressure capability to an existing blow-down facility at the Gas Dynamics Laboratory of the University of Illinois at Urbana-Champaign.

The findings in this report are not to be construed as an official Department of the Army position, unless so designated by other authorized documents.

1. FOREWORD

Work conducted under this grant pursued two long range objectives, namely:

- To investigate aerodynamic phenomena affecting the launch phase of in-tube launched rockets, and
- 2. To achieve direct compatability between in-house computer oriented (MICOM) efforts and the research conducted at the University of Illinois. This was to be achieved by acquisition of an HP 9830 system to be located in the Gas Dynamics Laboratory of the University of Illinois.

An immediate goal consisted of generating a comprehensive, yet well manageable analysis (including computer program development and typical performance documentation) of the launch and flight performance of a rocket with specified motor design and grain configuration.

Dr. H. H. Korst, Professor of Mechanical Engineering, acted as
Project Director; Dr. R. A. White, Professor of Mechanical Engineering,
supported by Mr. Dean H. Keal, was in charge of the facility development
in the Mechanical Engineering Laboratory in preparation of the experimental
effort which is to continue.

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Table 1 Parameters Affecting Internal Ballistic Performance

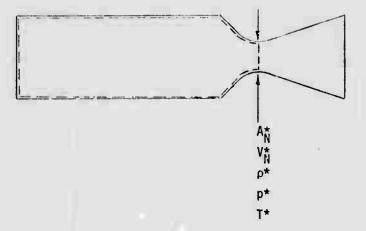


Figure la Control Volume for Evaluation of Mass

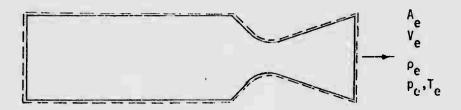


Figure 1b Control Volume for Evaluation of Thrust

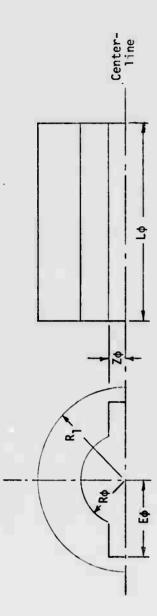
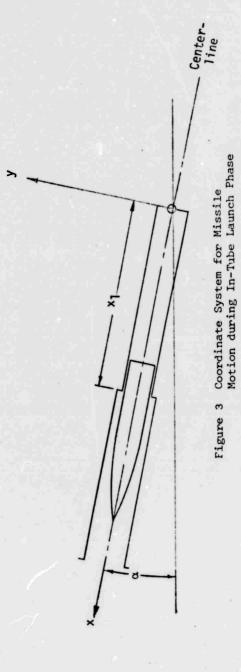


Figure 2 DOGBONE Grain (Geometry)



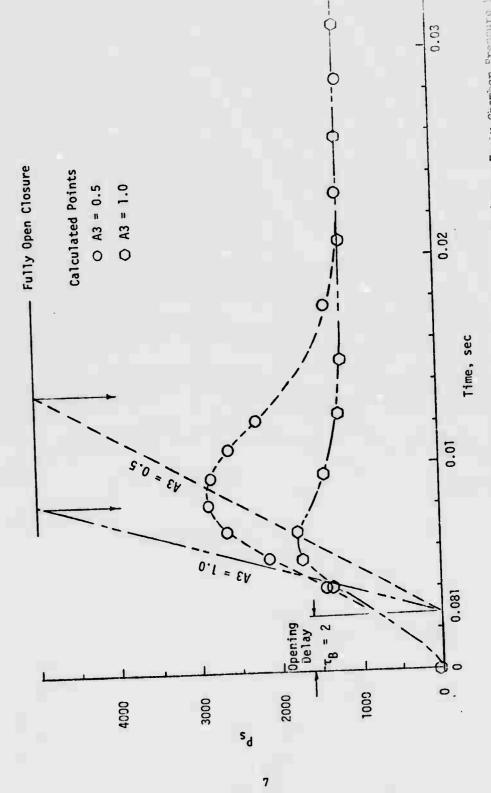


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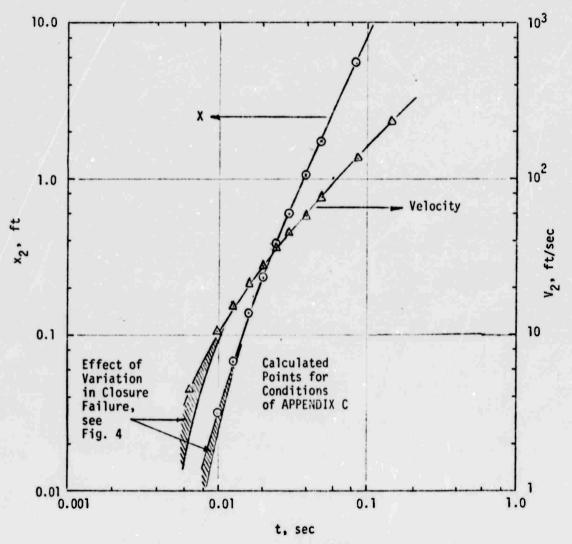


Figure 5 Effect of Closure Failure on Rocket Launch Speed

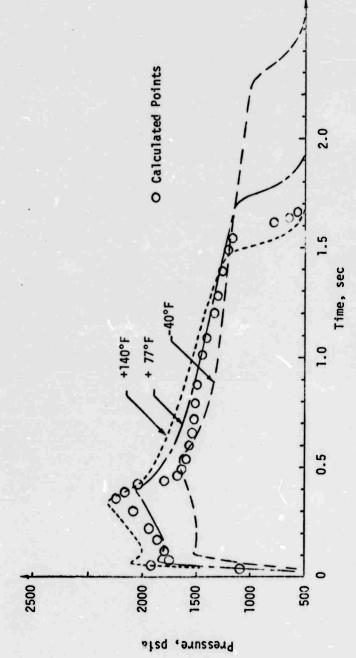


Figure 6 Chamber Pressure History for ARROW Rocket

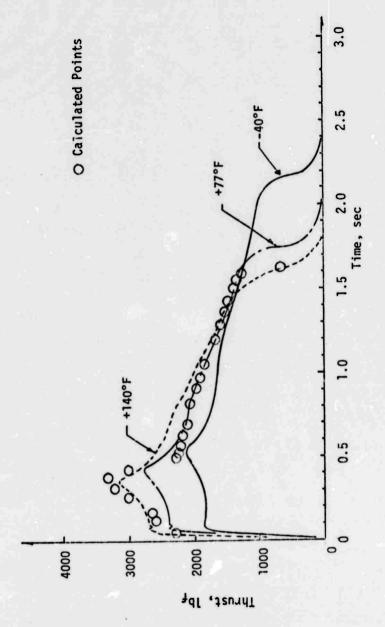


Figure 7 Thrust-Time History for ARROW Rocket

Table 1 Parameters Affecting Tuternal Ballistic Performance

Υ	Specific heat ratio, (-)
H	Molecular weight, lbm/mole
ToF	Flame temperature, °R
v _c	Co-volume, ft ³ /lb _{in}
v _s	Specific volume, ft ³ /lb _m
В	Burning coefficient, Eq. (17)
n	Burning exponent, (-), Eq. (17) R1, R ϕ , E ϕ , Z ϕ , L ϕ , grain dimensions, ft, see Fig. 2
Wo	Initial weight of round
L	Loading ratio, lbm/ft3
v _c	Chamber volume, ft ³
A**	Nozzle throat area, ft ²
A _e	Nozzle exit area, ft ²
P _{0,0}	Reference (atm) pressure, $1b_f/ft^2$
π _B	Burst pressure ratio, (-)
a	Effective linear opening coefficient, (-)

4. PERFORMANCE ANALYSIS OF A DOGLEG-GRAIN, TUBE LAUNCHED ROCKET AND AUXILIARY EFFORTS

During the launch phase, as grain burning is initiated, the propulsive nozzle remains initially closed to promote faster pressure buildup. Subsequently, the closure fails, starting at a prescribed burst pressure and allowing (critical) outflow through a (linearly) time variable effective throat. Neglecting fast transients, quasi-steady outflow conditions (together with the burning law and time varying grain surface) will determine the pressure and temperature history in the combustion chamber in which the gaseous phase of the propellant is assumed to be of uniform (stagnation) state. For fully established nozzle flow conditions (the initial shock system being expelled quicly), it is then possible to determine propulsive thrust, rocket acceleration velocity, and position as functions of time. Parameters considered for the present analysis are listed in Table 1. To investigate rocket-launcher interaction, it is then also of importance to investigate the effects of plume-wall interactions and the flow of propellant gases in the launch tube.

While analytical work on the propulsive rocket motion has been completed, certain experimental phases of rocket-gas interactions within the launch tube are still in progress and will be reported on at the conclusion of a follow-up effort under ARO-sponsorship.

4.1 CONTROL VOLUME ANALYSIS OF ROCKET MOTOR PERFORMANCE

Boundaries of the control volume are selected to include the entire missile including the propulsive nozzle (Fig. la); however, analysis of the chamber pressure as based on the conservation of mass utilizes the somewhat smaller volume terminated by the nozzle throat (Fig. lb). These, together with the assumption of a uniform state in the combustion chamber,

appear to be allowable simplifications [1].*

4.1.1 Conservation of Mass

With m $_{\rm S}$ and m $_{\rm g}$ denoting the mass of solid and gaseous phases of propellant in the chamber having the v-lume V $_{\rm C},$ one introduces the specific volumes to express

$$m_g v_g + m_S v_S = V_C$$
 (1)

and, by differentiation with respect to time, as $dv_s/dt = 0$,

$$v_{s} \frac{dm_{s}}{dt} + m_{g} \frac{dv}{dt} + v_{g} \frac{dm_{g}}{dt} = 0$$
 (2)

where $m_g = (V_c - V_s m_s)/v_g$ and

$$v_g = v_c + \frac{R T_o}{P_O}$$
 (3)

(Clausius-type gas).

With $T_o = F_{oF}$ and $v_c = constant$

$$\frac{dv_g}{dt} = -\frac{R T_{oF}}{p_U^2} \frac{dp_o}{dt}.$$
 (4)

The original mass of solid propellant at t = 0, m_{SO} , is used to form the dimensionless ratio

$$\mu = \frac{m_S}{m_{SO}} \tag{5}$$

and accounting for the outflow of gas from the control volume through the nozzle throat under choking conditions

$$\frac{dm_g}{dt} = -\frac{V_N^{\pm} A_N^{\pm}}{V_N^{\pm}} \frac{dm_g}{dt}$$
 (6)

For convenience, we now introduce an "ideal" acoustic reference velocity

^{*}Numbers in brackets refer to entries in REFERENCES.

$$C_{OR} = \sqrt{\gamma R g_{C} T_{OF}}$$
 (7)

and a characteristic time

$$t_o = \frac{V_c}{A_N^{+} C_{OR}}$$
 (8)

to define additional dimensionless variables,

$$\tau = \frac{t}{t_0} \tag{9}$$

$$\pi = \frac{P_O}{P_{O,O}} \tag{10}$$

$$\phi_{N} = \frac{v_{N}}{c_{OR}} \tag{11}$$

and parameters

$$c_1 = \frac{R T_{oF}}{P_{o,o} v_s}$$
 (12)

$$c_2 = \frac{v_c}{m_{so} v_s} \frac{1}{L v_s}$$
 (13)

Thus, the conservation of mass expressed by Eq. (2) attains the form

$$\frac{\mathrm{d}\mu}{\mathrm{d}\tau} \left[1 - \frac{\mathrm{v}_{\mathrm{c}}}{\mathrm{v}_{\mathrm{s}}} - \frac{\mathrm{c}_{1}}{\pi} \right] - \mathrm{c}_{2} \left[\frac{\mathrm{v}_{\mathrm{c}}}{\mathrm{v}_{\mathrm{s}}} + \frac{\mathrm{c}_{1}}{\pi} \right] \phi_{\mathrm{N}}^{\pm} \left[\frac{\mathrm{v}_{\mathrm{s}}}{\mathrm{v}_{\mathrm{N}}^{\pm}} \frac{\mathrm{A}_{\mathrm{N}}^{\pm}(\tau)}{\mathrm{A}_{\mathrm{N}}^{\pm}} \right] = \frac{\mathrm{c}_{2} - \mu}{\frac{\mathrm{v}_{\mathrm{c}}}{\mathrm{v}_{\mathrm{s}}} + \frac{\mathrm{c}_{1}}{\pi}} \, \mathrm{c}_{1} \, \frac{1}{\pi^{2}} \, \frac{\mathrm{d}\pi}{\mathrm{d}\tau}$$

$$(14)$$

For nozzle flow under choked conditions, by combining Eq. (3) with Bernoulli and energy equation, we obtain $A_N^{\pm}(\tau)$

$$\phi_{N}^{\pm} \frac{\mathbf{v}_{s}}{\mathbf{v}_{N}^{\pm}} \frac{A_{N}^{\pm}(\tau)}{A_{N}^{\pm}} = \begin{pmatrix} A_{N}^{\pm}(\tau) \\ A_{N}^{\pm} \end{pmatrix} \frac{\sqrt{\frac{2}{\gamma - 1}} \left(1 - \frac{T^{\pm}}{T_{of}}\right)}{\frac{\mathbf{v}_{c}}{\mathbf{v}_{s}} + \frac{C_{1}}{\pi \left(\frac{T^{\pm}}{T_{of}}\right)^{1/(\gamma - 1)}}$$

$$(15)$$

where

$$\frac{\underline{\mathbf{T}}^{*}}{\underline{\mathbf{T}}_{\text{of}}} = \left\{ \frac{\underline{\mathbf{Y}} - 1}{2} \left[1 + \frac{\mathbf{v}_{c}}{\mathbf{v}_{s}} \frac{\pi}{C_{1}} \left(\frac{\underline{\mathbf{T}}^{*}}{\underline{\mathbf{T}}_{\text{of}}} \right)^{1/(\gamma - 1)} \right] + 1 \right\}^{-1}$$
(16)

Attention is now given to the opening (failure) characteristics of the closure. The nozzle is originally closed by a diaphragm which leads to a delay in the outflow through the nozzle, also promoting the initial pressure buildup in the combustion chamber. The effect of burst pressure and opening coefficients on recoilless gun operation has been discussed in greater detail [1] where also experimental evidence is presented. We are retaining here this analytical concept by introducing a burst pressure ratio π_B at which the closure begins to fail at the time τ_B in such a way that it opens the throat area linearly with time at a rate determined by the opening coefficient, A_3 .

$$1 > \frac{A_{N}^{\ddagger}(\tau)}{A_{N}^{\ddagger}} = 0 \text{ for } \tau < \tau_{B}$$

$$A_{3}(\tau - \tau_{B}) \text{ for } \tau > \tau_{B}$$

Full opening is attained at time τ_1 where

$$\frac{A_N^{*}(\tau)}{A_N^{*}} = 1$$

after which time $A_N^{\alpha}(\tau)/A_N^{\dot{\alpha}} \equiv 1$. It is of interest to note that the failure characteristics of the closure exert some influence on the initial phases of the rocket launch especially on the peak chamber pressure, but seem to attenuate rather quickly after full nozzle aperture has been attained.

4.1.2 Propollant Co. Com: ion

Expressing the burning rate in conventional form by

$$r = B p^{n} (in/scc)$$
 (17)

accounting for the temperature dependency in P and the pressure dependency by the exponent n, we can determine the rate of propellant gas generation from

$$-\frac{\mathrm{dm}_{\mathrm{S}}}{\mathrm{dt}} = \frac{\mathrm{r} \, \mathrm{S}(\mathrm{t})}{\mathrm{v}_{\mathrm{S}} \, 12} \tag{18}$$

where S(t) is the time dependent burning surface of the grain, so that in dimensionless form

$$\frac{d\mu}{d\tau} = \frac{dm}{dt} \frac{d}{dm} \frac{d}{dm} \frac{dt}{d\tau} = \frac{S r t_o}{v_s m_{SO}} \frac{12}{12}$$
(19)

For the initial grain configuration given in Fig. 2, a solely time dependent burning rate (note that the pressure p is assumed to be uniform over the entire surface S), will produce changes in the grain geometry so that

$$R(\tau) = R_o + \frac{t_o}{12} \int_0^{\tau} r d\tau$$

$$E(\tau) = E_o + \frac{t_o}{12} \int_0^{\tau} r d\tau$$

$$Z(\tau) = Z_o + \frac{t_o}{12} \int_0^{\tau} r d\tau$$
(20)

The burning surface $S(\tau)$ can now be expressed as follows:

(i) For
$$Z^2 + E^2 \le R1^2$$
, one obtains
$$S = 4L\{R(90 - \theta) \pi/180 + E + Z - K \cos \theta\}$$
 (21.1) where $\theta = \sin^{-1}(Z/R)$.

(ii)
$$S = 4L \left[R(90 - \theta) \pi/180 + \delta R_1 \sqrt{1 - (E/R1)^2} + R_1 \sqrt{1 - (Z/R_1)} - R \cos \theta \right]$$
where $\delta = 1$ for $E < R1$ (21.2)
$$= 0 \text{ for } E > R1$$
 (21.3)

Also, for the grain configurations, Fig. 2,

$$\mathbf{m}_{so} \mathbf{v}_{s} = \mathbf{L} \mathbf{R}_{1}^{2} \left\{ \pi \left[1 - \left(\frac{\mathbf{R}_{o}}{\mathbf{R}1} \right)^{2} \right] \left[1 - \frac{\theta_{o}}{90} \right] + \left(\frac{\mathbf{R}_{o}}{\mathbf{R}_{1}} \right)^{2} \sin 2\theta_{o} - 4\mathbf{R}_{o} \mathbf{E}_{o} \sin \theta_{o} \right\}$$
(22)

where $\theta_0 = \sin^{-1} (Z_0/R_0)$

4.1.3 Chamber Pressure-Time History

Noting that

$$\frac{r t_o}{12} = \frac{B p_{o,o}^n \pi^n v_c}{12 A_N^2 c_{oR}}$$
 (23)

It is possible to calculate the pressure-time history in the combustion chamber by integrating Eq. (14) together with Eq. (19) once the grain geometry (Eqs. (21) and (22)) and the burning law (Eq. (23)) are specified.*

4.2 ROCKET DYNAMICS AND KINEMATICS

4.2.1 Initial Phase

Considering the rocket motion here as resulting from nozzle thrust and the gravitational acceleration only and using a Cartesian system of coordinates aligned with the launch tube axis (see Fig. 3) which

^{*}A summary of parameters introduced into the present analysis is given in Table 1.

is inclined to the horizontal by the angle a, one obtains

$$\frac{dV_{x}}{dt} \frac{W}{g_{c}} = T - W \frac{g}{g_{c}} \sin \alpha \qquad (24)$$

and

$$\frac{dV_y}{dt} \frac{W}{g_c} = -W \frac{g}{g_c} \cos \alpha \tag{25}$$

In restricting ourselves to these expressions, we have, for the time being, neglected all aerodynamic forces acting on the rocket. We note, however, that the lateral motion of the rocket which begins after the rocket has already moved the distance \mathbf{x}_1 at time \mathbf{t}_1 in the tube, may well give rise to unbalanced aerodynamic interference forces which can lead to pitching during launch.

For the initial phases of the lateral motion, we may neglect changes of the rocket mass W and, using an average thrust \overline{T} , integrate Eq. (25) twice to yield, as y = 0 and $V_y = 0$ at $x = x_1$, $t = t_1$.

$$y = -\frac{\cos \alpha x_1}{\frac{T}{W_c} g_c} - \sin \alpha \left(\frac{t}{t_1} - 1\right)^2$$
 (26)

The time t_1 where the distance x_1 has been negotiated is found from the integration of Eq. (24) for which $V_x = 0$ at t = 0 and

$$x = g \frac{t^2}{2} \left(\frac{\overline{T} g_c}{W g} - \sin \alpha \right)$$
 (27)

so that

$$t_1 = \frac{2x_1}{\frac{T}{W} g_c} - g \sin \alpha$$
 (28)

4.2.2 Treet of Variable Thrust and Mass

If the fall trajectory is to be followed, it will be necessary to introduce eredy, ric lift and drag forces and to consider also the time variations of rocket thrust and mass. Accounting for the latter two only, one realizes the results of pressure chamber and propellant mass ratio histories as they result from Section 4.1.

While Fig. 4 illustrates the influence of closure failure on the early chamber pressure history, Fig. 5 shows that the effect on launch velocity is indeed very small.

The full chamber pressure history, as calculated by the present program (for instantaneous closure failure), is compared to available ARROW data in Fig. 6.

Accounting for the variable rocket mass by

$$W = W_{o} \left[1 - \frac{ms_{o}}{W_{o}} (1 - \mu) \right]$$
 (29)

and determining the thrust force on the basis of the control volume shown in Fig. la, the time-dependent thrust force is given (in 1b_) by

$$T = P_{o,o} A_{N}^{\pm} 1^{44} \left[\frac{2\gamma}{\gamma - 1} \frac{\sqrt{\left(1 - \frac{T_{e}}{T_{oF}}\right) \left(1 - \frac{T_{n}^{\pm}}{T_{oF}}\right)}}{\frac{v_{c}}{v_{s}} \frac{1}{C_{1}} + \frac{1}{\pi} \frac{T_{n}^{\pm}}{T_{oF}}} + \frac{A_{e}}{A_{N}^{\pm}} \left[\pi \left(\frac{T_{e}}{T_{oF}}\right)^{\gamma/(\gamma - 1)} - 1 \right] \right]$$
(30)

Again, the thrust force calculated with the present program is compared with data provided for the ARROW rocket (see Fig. 7). The temperature ratio $T_{\rm e}/T_{\rm oF}$ is found after solving Eq. (16) for $T^{\pm}/T_{\rm oF}$ by iteration, from

$$\frac{T_{e}}{T_{o\Gamma}} = \left\{ \left[\frac{\pi}{C_{1}} \frac{v_{e}}{v_{s}} + \left(\frac{T^{\pm}}{T_{o\Gamma}} \right)^{1/(1-\gamma)} \right] \sqrt{\frac{1 - T_{e}/\Gamma_{o\Gamma}}{(1 - T^{\pm}/T_{o\Gamma})}} \frac{\Lambda_{e}}{\Lambda_{\parallel}^{\pm}} - \frac{\pi}{C_{1}} \frac{v_{e}}{v_{s}} \right\}$$
(31)

which also a quires iteration.

4.3 COMPUTER PROGRAM DEVELOPMENT

As already pointed out in Section 1, the broader objectives of this grant called for computer program compatibility between MICOM and the supporting effects at the University of Illinois. Consequently, after acquisition of the HP 9830 system, a considerable number of computer programs have been generated which are now operational at both locations.

Some representative examples are cited in the following

- 4.3.1 Internal Ballisties Program
 - 4.3.1.1 Bomb, Gun, and Recoilless Gun Performance Analysis
 (Documented earlier [1].)
 - 4.3.1.2 Adaptations of such Programs to Deal with Gun-Launched Rockets [2]
- 4.3.2 Plume and Slipstream Boundary Analysis based on the Method of Successive Centered Expansions [3,4]
- 4.3.3 Base Pressure Analysis for Unpowered Flight of Rockets and Projectiles [5,6]
- 4.3.4 Viseous Jet Mixing and Boundary Layer Programs in Support of Drag Evaluations for Missiles [7]
- 4.3.5 Shock Interaction Programs Applicable to Muzzle Break Blast-wave Propagation and Reflections [1]
- 4.3.6 Comprehensive Program for the DOGBONE Grain (ARROW) Rocket Performance (based on the analysis of this report)

4.3.6.1 Program Printout

A printout of the program for the HP 9830 computer is given in APPENDIX A.

4.3.6.2 User's Instructions

User's instructions are given in APPENDIX B with all input quantities defined (also, see Table 1).

4.3.6.3 Program Output ("PRINT ALL" Mode)

Program output in the "PRINT ALL" mode of the comuter is given in APPENDIX C for selected input data as listed. It must be noted that the DOGLEG GRAIN geometry (R1, R ϕ , E ϕ , Z ϕ) and nozzle area ratio $\Lambda_{\rm e}/\Lambda_{\rm N}^{\rm ii}$ = A5 are not entered from the keyboard but are READ by the program from the DATA line 120. Changes in these parameters have thus to be made (if so desired) by "FETCH 120, EXECUTE" and by altering the values in this line (conveniently done by using the editing features of the HP 9830).

4.4.1 Facility Development

The blow-down wind tunnel facility and auxiliary air supplies of the Mechanical Engineering Laboratory, Department of Mechanical and Industrial Engineering, University of Illinois at Urbana-Champaign, have been modified to allow modeling of the tube launch system under quasisteady operating conditions. The pressure distribution within the launch tube due to plume-tube wall interactions and vehicle eccentricity is the primary objective.

4.4.1.1 Air Supply Hodifications

sary to modify the existing auxiliary air supply system to accept the higher pressures needed for simulating the rocket jet plume. This has been accomplished by the installation of new high pressure storage (approximately 450 ft³, maximum working pressure 1800 psig) and high pressure piping in compliance with OSHA regulations. A two-stage compressor allowing pumping of the system to either 250 psig or 500 psig levels is currently being utilized.

4.4.1.2 Model Constructed

A one-half scale geometrically similar model has been constructed of both the afterbody and the launch tube system. The model consists of an interchangeable nozzle section installed in the end of a section of schedule 80 high pressure pipe with nominal outside diameter of 1.90 inches. The launch tube is simulated by two lengths of thin wall brass tubing with nominal inside diameters of 2.0 and 2.25 inches.

4.4.2 Preliminary Experiments Conducted

A series of calibration checks with the nozzle configuration selected were carried out to determine the performance characteristics of the modified high pressure system and its control valve. The results of the preliminary tests indicated that maximum stagnation pressures of only 180 psig are presently reached with tank pressures of approximately 450 psig. This is below the level desired for proper plume simulation within the launch tube. It appears that the flow level capability of the control valve is too close to the nozzle mass flow required or that a sonic throat

is occurring within the pipe system restricting the available pressure to the nozzle. Additional tests are being planned to determine and subsequently improve the range of system performance. A reduction in model scale to approximately 0.25 may be required.

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Computer Program "BOGLE: C:AIN MCCKET" Printo: L for HP 9830 N-K (no LCCC)

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1000
to million to be a to a control of
THE DEF FRECKE HORE ! WITH A MARCH
15 [10]
20 UFG
25 DISP "GAMMA" )
SO INCULT
35 JUSP WOLEDBERG MAIGHTENESS !
40 HIPUL II
45 Insk TREAME TENL, Dec 25%;
50 THEUT TO
$5 CO=(K*32.1 4*1545*TOZH) P0.5
60 DISP "COVOLUME PROP. VOH";
65 IMPUT VI
70 DISP "CHAMBER VOL. V=' 3
75 IMPUT V
80 DISP "SPEU. VOL. SOLIDO VO=";
85 THPUT V2
98 JUSP "LOAD RATIO:L=";
95 IMPUT L
188 HISP "INITIAL MACS: MISSILE")
105 IMPUT NO
118 C2=1/(L*V2)
115 READ RI-RA-E0,20-A5:LA
120 DATA 0.14167:0.057297:0.115:0.015625:3.7603:2.5958
125 PRINT "R1='R1;"R0:"R0;'E0="R0;"Z0="Z0;"NEZ8/-"A5
130 DISP "BURH. COEFF. B=";
135 INPUT B
140 DISP "BURN. EXP.N=";
145 IMPUT N
150 DISP " BURST PRESS. RATIO =";
155 INPUT B2
160 DISP "AREA HOZZLE: A1=";
165 IMPUT AL
170 DISP "P0:REF="3
175 THPUT PO
180 T9=V:(A1x00)
185 C1≈1545÷T0∴(144×P0×V2×H)
190 DISP "OPENDIG COEFF.A="4
195 INPUT A3
200 DISP "TIME INCREM. (1)="3
205 THPUT T
210 PRINT "GANHA="K)"M="h;"T0="Y0;"C0V0L VC="V1; CHAnb.V0L.V="V
215 PRINT "LOOD RATIO L≃'L3"SPEC. VOC.SDL.VS: "V2: 'INITIAL MASS, MISS. 'NO
220 PRINT "BURN COEF.B="B) "BURN EXP.N="N
225 PRIM 'BURSI PLATE PRESS.="BX
          "AREA HOZZLE A1="A1;"KET PRESS. PG="FO;"OPEN. LOEFF.HL"KA="HB
200 PRIME
235 PRINT
240 V3=V1/V2
245 D0-FNS(70/R0)
250 B1=80(P0)4H
255 D3:PI*(1-(R6/R1)+2*(1-D6/90))+/R6/R1)+2*(Id(2+D6)-4/E6*R6*SHD6-F1+2
```

269 Di=-Bi*T9/(3*Ri*F3)

```
BER EFLO
ans Hara
SEL PI N
0 10 20:0
325 194
330 63-1
335 P9-P1
348 D=FN (25)
355 A=1
360 0≈(((1+V<sub>5</sub>∞ersiAra + i ≥ 1) (4 milion ki 1) (2+1) (1 (1)
365 IF ABS(A (D(o, A, E))) | 104 534
370 A=0
375 GUTO 366
380 H=(2)(1+0) di-c bare Vs diedrift-KovEDY
385 1F G1-1 TREE 400
398 TP B20P1 THEN A10
395 IF GP=1 THER 625
488 N4:A3:(11 18)
405 JF 8401 BHEN 645
418 H2≃M1+89÷T
415 Y1=1 (P9*) 98 VU (0 = 11) / (2 - (1) - 112) /(2)
428 Y2=B9%(1.V3.C1/19)
425 Y3=+02*(V3+01/P40
430 Y4=Y0+04*H
435 fr2=fr1+Y1x,Y2+Y4)
446 P9=(P21P1)/2
445 03=0.55
450 01=((F9*V3:01+04-12/1 F>))*((1-02)/(1-0))*0.5*05-P9*V3/C1/f(1 k)
455 IF MBS(01-02):0.231 THEH 478
460 02=01
465 G010 450
470 F1=2+K+(<1-01)+(1-0))43.5/((K-1+K(V3)C1+U*(1+1-K+)/P9-)
475 f2=A5< (P9+()1*()+((<-1))+1)
488 F3=P0*N1«(F1+F.')*144
485 | 12: | 11:64×F3:T7T0:22,124:(MO: 1-(L0:K112:H3:(V2:H6:O::) 0::1 (M2:HD ::2/):
490 X2:X1:(N3:N1)/1:019/2
495 IF G3#1 THEN 500
500 GOTO 515
505 G3+0
518 6010 345
515 II A4#1 THEN 525
520 / I#10 THEN 555
538 PRINT "ASKTI2 H2 - A43772KSEC.="G_PH ( → 10
535 PRINT "THRUST (LBF)="13;"VELOCITY (LPG)="48;")2(FI)="48
540 PRINT
```

```
1000 1-111
506 Pl-P.
5655 hts R2
570 T1=T101
575 DP-19-84 Fe 1 1
560 R RED
585 FFR402
590 Jr 2403
595 MI-NO
600 X1-X2
605 IF P941 THEIL : 55
610 1F M2:0 CHER 300
615 111-8
620 6010 330
6.5 62-0
638 T84T1
635 61-1
648 6010 400
645 H4=1
650 GOTO 410
695 F=1
660 IF RIVE THEN 670
665 F=8
670 B8=(R*(99-D)*FID:150-R1 -F-CO5D P1+Fadd 1 805-1 CE/R1/12.):
675 89=D1 P9TH+(E8+SCP(1-+Z-F1)*2))
688 GU10 355
685 STOP
690 EHD
```

APPENDIX B

User's Instructions for "DOGLEG GRAIM ROCKET" Program

ARROW-BALLISTIC PROGRAM

HP 9830 (no ROMS needed) 690/1428 "DOGBONE GRAIN ROCKET" load

Inp	ut (asked for by program)	Symbol	
1.	Gamma (propellant), 1	K	
2.	Molecular weight propellant, 1	М	
3.	Flame temperature, propellant, °R	ТØ	
4.	Co-volume, propellant gas, ft ³ /lb _m	Vl	
5.	Combustion chamber, volume, ft ³	ν	
6.	Specific volume, solid propellant, ft^3/lb_m	V2	
7.	Load ratio, lbm/ft3	L	
8.	Initial mass, missile and propellant, 1b m	MØ	
	(location 120 generates (can be altered) Data R1, R ϕ , E ϕ , $7\phi, A5 = A_e/A_N^{\circ} and L\emptyset related to the DOGBONE geometry and nozzle configuration as shown in Figs. 1a, 1b, and 2)$		
9.	Burning coefficient, in./see	В	
10.	Burning law, exponent, 1	N	
11.	Burst pressure ratio, 1	B2	
12.	Area nozzle (throat A*), ft ²	Al	
13.	Reference pressure PØ, psia	Pø	
14.	Opening coefficient (burst plate), 1	АЗ	
15.	Time inerement ($\Delta \tau$), 1	Т	

NoTE: It is a still to use T = 0.5 for accuracy and numerical stability.

Compute proper, after normle throat has fully opened (after fast transients of burnt plate of ming have been recorded) skips intermediate PRINTOUTS

(retaining T), gives data for only every tenth time increment. This will generate one set of data every 100 seconds (approximate real computer time).

Program stops when combustion chamber pressure reaches PØ.

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111.
Jacobson Land, for
Motor Claum In 1 and a second
TERME TERMS INC. . . .
CUMPLEM PIRE DO DE
HIMMAR Man Vit Line
SPEC. VOL. SOLIDS VS 7.0.. .
FORD RHYICOLE 4.58
INTIDIC METAL AND SHEET IN
BEAR = 3.76031-404
BURN COLFF. R=9.8250
EURN. ENF. N=1.46
 BURST PASS. RATIO : "TH
AREA NOZZEŁA ALY JOSE BUZZE
1 8: REF-714: 4
OPENING COEFF.H=91
TIME INCHEM : 19=7.1
CAMMA: 1.1835EH00 M- 3.665 0 40 TH- 5 45,70 40 COVOL YO
INTIDAL MASO: MISS = 4.709(F+01
BURN CORF.D= 0.5060E 82 BURN EAR.h- 4.5080E-01
EURST PLATE FRESS = 1.00005+01
AREA NOZNUE A1= 8.9908E OF THE PRESS FOR 1.4800F HUI
OPEN. COUFF. PLPNA: 1.0000F For
B2(Ti)/HC=
 THRUST (ERF)≈ 2.8354E+01 | VELOCITY (FMCC= 0.5686E+U0
                                                                                                                      112(F1)=
 0.0000E+60
TAU1: 1.00006-01 FRESCA 1.9956E400 MH2: 9.9979E-01 M2(F1):82=
  0.8809E 00 T2/SF( = 1.8879E 00
32(FT)=
 0.0000000
TRUIT 2.00006-01 | FRESS= 5.67056+62 | NUCT | 20506-01 | H2CTI / M25.
 0.0040E+00 T2<°EC>= 1.6714E+03
THRUST (LOF)- 5.1717E+02 VELOCITY (FES)- 1.8000E+00
                                                                                                                      82 (FT)=
  0.00000000
THEO: 3.8008E 01 FEESS 5.0917E/02 MURE 9.9937E-01
                                                                                                                        02071J 624
 1,6000E-01 12(SEC)= 7.1757(-82
 THE USA (FRE)= 9.98/78/53 - VELOCIAY (FRO)A 0.2849E-88
                                                                                                                      RESERVE
  8.93075-06
 TAULT SECRET OF THE PROPERTY O
  2.60098-01 [2 :01:: 1.71974 [6]
 THRUST CERTA LARGER OLD PRINCIPLY CEPTA-1, BUT IN XXVII o
  5.30490:00
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1 the set of 4.41011 (1) 1001 7.850 0 0 10 1 970 1 10 1. 3010 00 40 1. 30 10.660 01 0 - 1. 50 0 2 1.150 05 10.660 01 01 1 1.50 03 70 00 1.1 10 0 1 11 10 0 72 0 72 0 1 9. 19. (5) - (1) Martin 8.00006-01 (1885) 1 4 38 (40) MU2= 1.57616 80 (6.1) 10 (6.2) 6.8888F 01 [2.860)= 4.8984F-03 THPUS! (LBY): .85088.408 YELOCI: (FPS) 1.0090E+00 X20ETO= 1.09658-03 TAUX = 9.00001-03 Prison at on the 102 of 7281-04 Addit At 7.0000E-01 70 SE100 5.4: E 0 2,95566-63 4.4748E 03 TAUT= 1.1000E FOR PRESS 1.0450EF03 NO. - 9.9055E NO. - 62.717 MCS 6.62446-83 92(11) Hus-1.0000E+00 12KSE()= 1.1947E-02 THRUST (LPC)= 2,47980:03 | VELOCITY (FPS = 1.13568+01 ha FIDE 5.88386-92 TARDS 3.1000E+06 FRESS: 1.1739E+03 HUZ- 1. 8670 61 P2(710/104 1.8000E+00 TP(SEC)= 1.7406E-02 THRUUT (LEF)= 2.1536E003 VELOCITY (FPST= 2.1117E041 PARE DE 1.60998 01 1801: 4.1000[+00 FRESH 1.1672[+03 M92= 5.8405E-0] H21.11 H28 1.8000E+00 1.2(SEC)= 2.2845E-02 HOOF DA THRUNE (LEF)= 4.00455100 YELOUDTY (FPS): 3.16736401 3, 1000E 01

1888 5.160956×60 PSESS- 1.17168+88 782= 9.61218 01 82×11×363= 1.080064+60 12×100.0× 7.92836 00 18885 6.04 2.-453€ 60 7810075 1850- 4.82896+81 82×866

5.05.28-01

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  Allia II, Ibraha ah
 Taking the programmer of the first transfer of the contract of
   1.3755L+80
TRUI 9.100(1:00 ).re -- 1241 kmrs - 800 - 120 % o - 120 1.020(6:00 - 720 %) o - 120 4-e 50
1.020(6:00 - 720 %) (-- 120 4-e 50 %)
TNRO 3. (LOF): 2.40 (-- 130 MELOCKIT KER- 120 220 4-0) - 40 4
  1.7cH4E +88
1.000 Except Ta End From Holf end
THRUS (LBF) 2.42062/82 VLEOCTT (FIST) 1.11/1/31 2010-
   3.1939E+08
TAUD: 1.1100E+01 FRESU 1.2125.100 AUG-0.5-06-01 HP315 AUG-1.0000E+00 TL45E.3-6-3621E-02
 THRUST (LOF)= .4.63F+03 /LLOCITY (LES:= 9.01267+01 ).1(F1)=
   2.67605+00
1.0000F+60 TP(SF1)= 0.0360F 0.0
THRUST (LBF): 2.45-25/02 VELO IT? (FPS) 1. EDME 050
                                                                                                                                                                                                                                   DIAMETER .
    3.20715+00
1.0000E+00 12:SEC> 7.1,90E 02
11.14FT)=
  3.7871E+08
4.418bEHBU
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